

# Jets and droplets from bursting bubbles

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## Introduction

 Started using Gerris for my PhD work here in Princeton (Craig Arnold's lab) in 2012



• Started postdoc at BU (Jacy Bird's lab) August 2016



Rising bubbles beneath inclined walls

Self-similar breakup of liquid cones

Jet drops from bursting bubbles

#### Jet drops from bursting bubbles



## Jet drops from bursting bubbles



## Motivation

- Atmospheric science
  - Sea spray aerosol particles act as cloud condensation nuclei, scatter radiation
    - Still significant uncertainties in climate forcing by aerosols



Richter & Veron 2016







## Dynamic similitude

- Nondimensionalization:
  - Length:  ${\cal R}$

– Time: 
$$\tau \equiv \sqrt{\rho R^3/\gamma}$$

- Neglect gravity (valid for  ${\rm Bo} \lesssim 0.01)$ (Bo  $\equiv \rho g R^2/\gamma \rightarrow R \lesssim 0.3~{\rm mm}$ )
- Only dimensionless parameter: Laplace number  ${\rm La} \equiv \rho \gamma R/\mu^2$ 
  - Note:  $La = 1/Oh^2$



#### - Increasing $\mu$ equivalent to decreasing R

Density, surface tension, viscosity:  $\rho \qquad \gamma \qquad \mu$ 

## Bubble bursting experiments

- Use glycerol-water solutions of varying concentrations to change viscosity, keeping  $\,R\approx 200~\mu{\rm m}$ 



## Bubble bursting experiments



 Non-monotonic relationship between size of the top jet drop and Laplace number

## **Bubble bursting experiments**



 Non-monotonic relationship between size of the top jet drop and Laplace number

## **Previous simulations**

- Duchemin et al. 2002 also observed nonmonotonic relationship in simulations
  - Limited resolution
  - Gap near minimum



# Numerical simulations

- Axisymmetric simulations run in Gerris
  - Adaptive mesh refinement in regions of high curvature, vorticity: max level 14
  - Minimum cell size =  $2.4 \times 10^{-4} R$
- Initialized as spherical bubble with popped cap
  - Neglect gravity
- Vary La, fixing  $\rho_g/\rho = 1.2 \times 10^{-3}$  $\mu_g/\mu = 0.018$



#### Validation



• Air bubble in water with  $R = 210 \ \mu m$ (La = 18000)

## Simulation results



- Define inversion time  $t_0$  as time when velocity of interface at center is maximum
  - Time of pinch-off  $t_p$  also labelled
- Same non-monotonic relationship

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- Decompose drop size into
  - Shape factor  $r^* \equiv r_d (\rho/\gamma)^{1/3} (t_p t_0)^{-2/3}$
  - Jet growth time  $t^* \equiv (t_p t_0)/ au$

• Then 
$$r_d/R = r^*(\text{La})(t^*(\text{La}))^{2/3}$$



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- As La decreases below 1200:
  - Viscosity delays pinch-off
  - Drop size increases with  $t^*$



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• As La increases above 1200:

- Less focusing of cavity with undamped capillary waves (Ghabache et al., 2014)
- Drop size increases with  $r^st$







#### Size variations due to escape from pinch-off



Hoeppfner & Paré (2013), Recoil of a liquid filament: escape from pinch-off through creation of a vortex ring

#### Dimensional plot: Jet drop radius vs. bubble radius



- Seawater viscosity varies by almost a factor of 3 from 0 °C to 40 °C → strong La dependence on temperature
  - Drop size increases with temperature for  $~R\gtrsim50~\mu{
    m m}$
  - Jet drops as small as 200 nm predicted in tropical waters

## Conclusions

- Non-monotonic size relationship between bubble and top jet drop observed
- Decomposing self-similar jet growth into shape and time components can capture non-monotonic behavior





Jet drop sizes 1 mm
 predicted
 significantly
 smaller than 2 10 µm
 10% rule and 1 µm
 temperature dependent



## Acknowledgments

- National Science Foundation Grant No. 1351466
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#### Self-similar scaling: La=610



#### Self-similar scaling: La=1700



#### Self-similar scaling: La=7200



#### Gravity effects: Bo~0.2



#### Gravity effects: Bo~1

